

Thermodynamic properties of correlated fermions in lattices with spin-dependent disorder

Karol Makuch

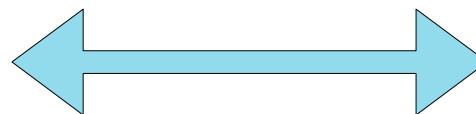
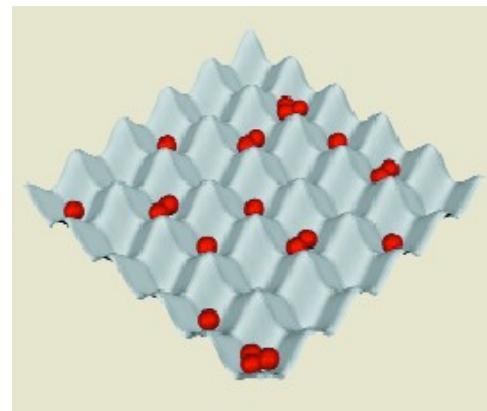


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D. Vollhardt (*University of Augsburg*)
arXiv:1302.3395

Warszawa, 21.03.2013

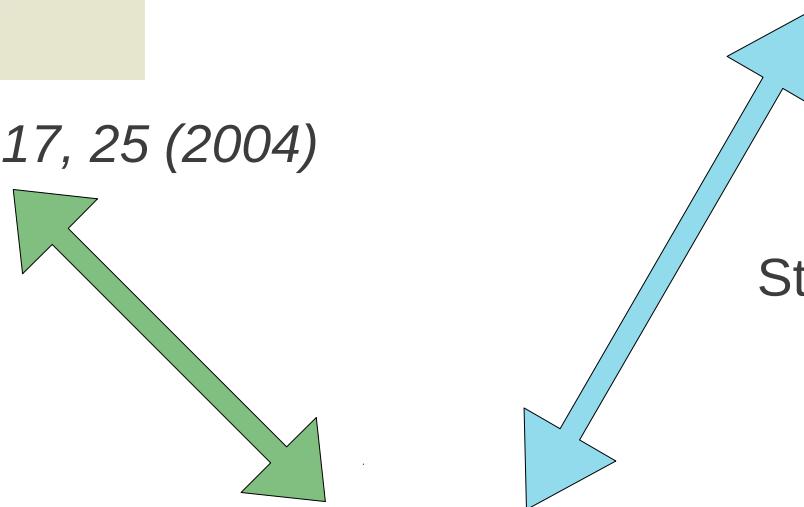
Introduction

Neutral particles in optical lattice



Interacting electrons in
materials

I. Bloch, Phys. World 17, 25 (2004)



Strongly correlated materials

Hubbard model (1963):

Electron correlations in narrow energy bands

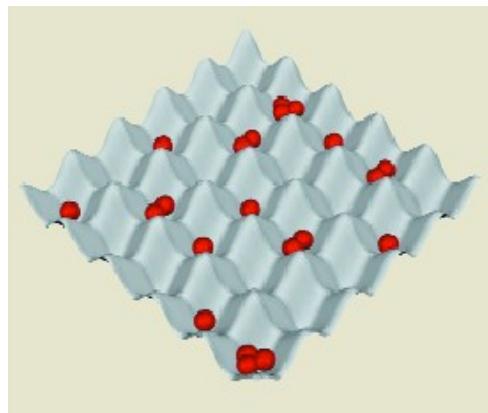
BY J. HUBBARD

that a possible approximation is to neglect all the integrals (8) apart from I . If this approximation, the validity of which is discussed in greater detail below, is made, then the Hamiltonian of (6) becomes

$$H = \sum_{i,j} \sum_{\sigma} T_{ij} c_{i\sigma}^{\dagger} c_{j\sigma} + \frac{1}{2} I \sum_{i,\sigma} n_{i\sigma} n_{i,-\sigma} - I \sum_{i,\sigma} \nu_{ii} n_{i\sigma}, \quad (10)$$

Introduction

Neutral particles in optical lattice



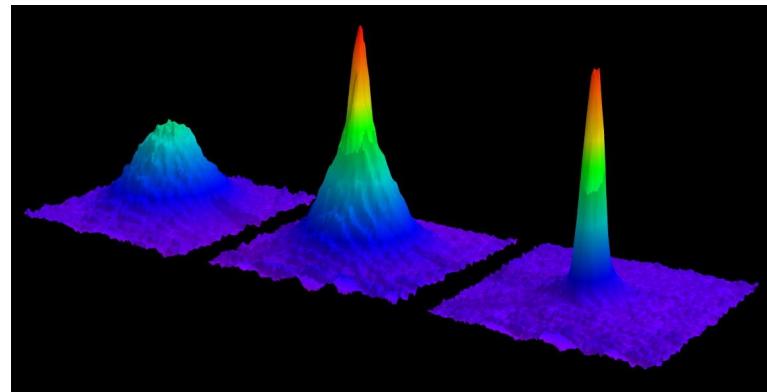
perfect crystal

very good realization of Hubbard model

Hubbard model not solved - optical lattice as quantum simulator
Searching for new quantum phases

Introduction

(1995) BEC: Eric A. Cornell, Wolfgang Ketterle, Carl E. Wieman



Introduction

Strongly correlated phase transition for the first time in optical lattice, bosons:

Quantum phase transition from a superfluid to a Mott insulator in a gas of ultracold atoms

Markus Greiner*, Olaf Mandel*, Tilman Esslinger†, Theodor W. Hänsch* & Immanuel Bloch*

NATURE | VOL 415 | 3 JANUARY 2002 | www.nature.com

Fermions:

A Mott insulator of fermionic atoms in an optical lattice

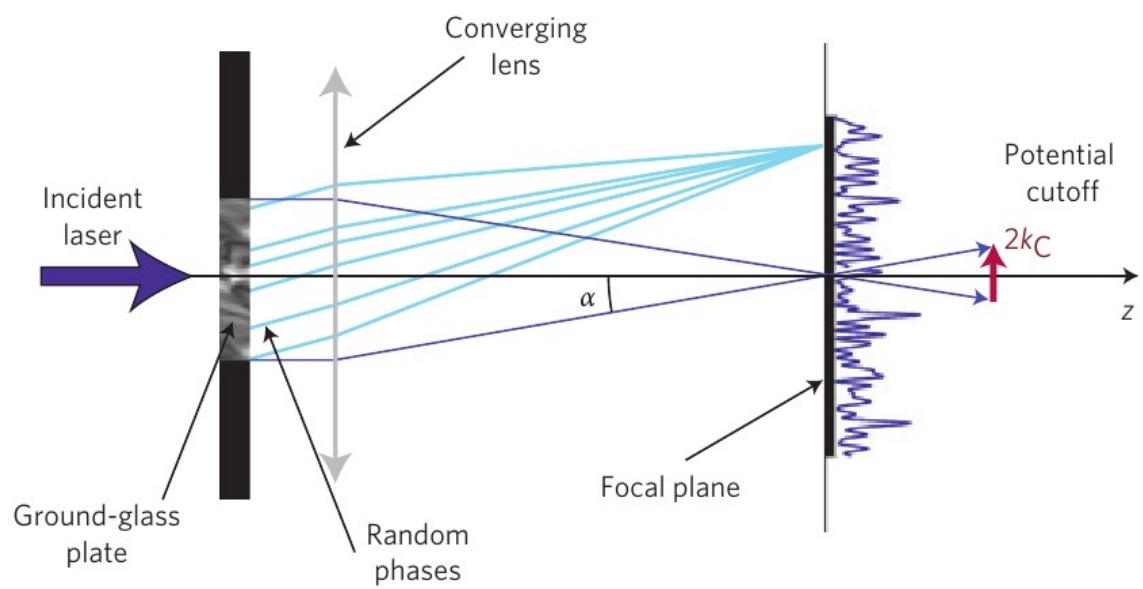
Robert Jördens¹*, Niels Strohmaier¹*, Kenneth Günter^{1,2}, Henning Moritz¹ & Tilman Esslinger¹

NATURE | Vol 455 | 11 September 2008

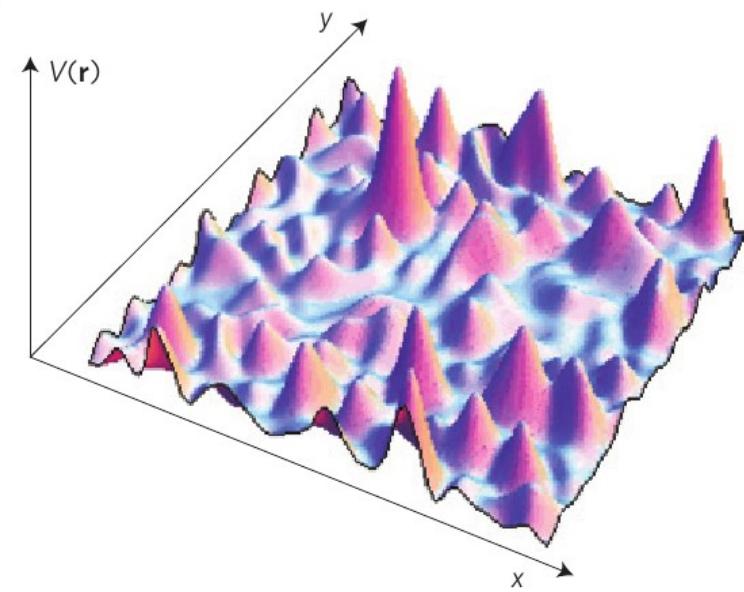
Introduction

Optical lattice with disorder

a



b



L. Sanchez-Palencia, M. Lewenstein, Nature Phys. 6, 87 (2010)

Motivations and aims of our work

Apart from disorder in optical lattices:

Spin-dependent lattices first proposed:

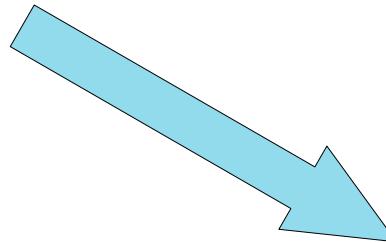
*D. Jaksch, H. Briegel, J. Cirac, C. Gardiner, and P. Zoller,
Phys. Rev. Lett. 82, 1975 (1999)*

Spin-dependent lattices first implemented:

*O. Mandel, M. Greiner, A. Widera, T. Rom, T. Hänsch, and I. Bloch,
Nature (London) 425, 937 (2003)*

P. Soltan-Panahi, J. Struck, P. Hauke, A. Bick, W. Plenkers, G. Meineke, C. Becker, P. Windpassinger, M. Lewenstein, and K. Sengstock, Nat. Phys. 7, 434 (2011)

- disorder in lattice realized
- spin-dependent lattice realized



Spin-dependent disorder in optical lattice
possible to realize
(beyond standard solid state physics)

Comprehensive thermodynamics?

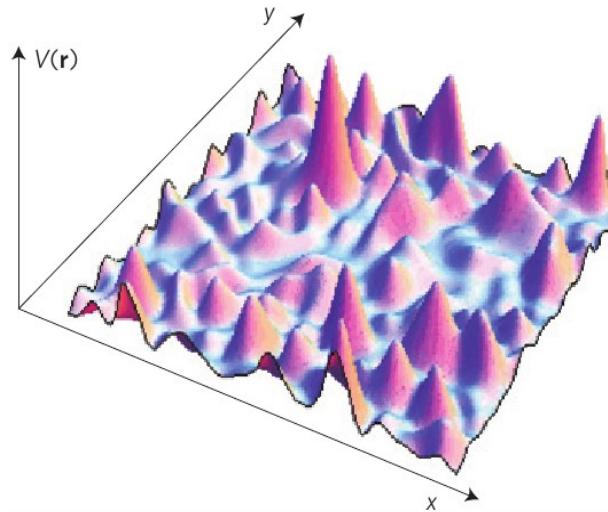
Hubbard model with spin-dependent disorder

$$H = t \sum_{(i,j)} \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} \hat{n}_{i\sigma} + \sum_i U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

$$\hat{n}_{i\sigma} \equiv \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

Disorder in the model:

- comes through diagonal terms
- quenched disorder



Similar model: *R. Nanguneri, M. Jiang, T. Cary, G.G. Batrouni, and R.T. Scalettar, Phys. Rev. B 85, 134506 (2012)*

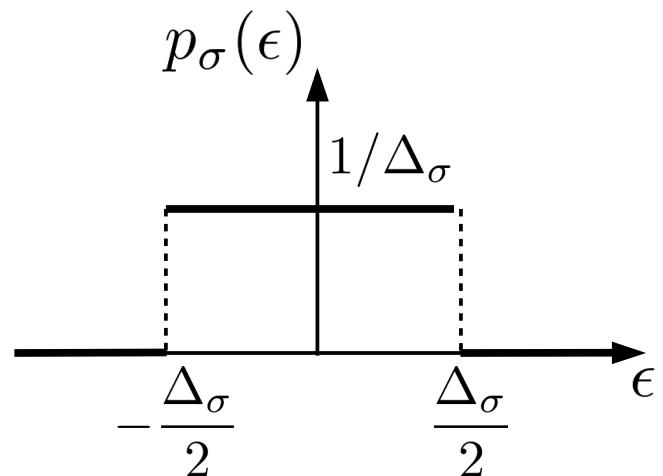
but: $U < 0$, Bogolubov de Gennes mean field theory

More about disorder in the model

$$H = \dots + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} \hat{n}_{i\sigma} + \dots$$

We assume **uncorrelated rectangular probability distribution function**:

$$P(\epsilon_{1\uparrow}, \epsilon_{1\downarrow}, \dots) = \prod_i p_{\uparrow}(\epsilon_{i\uparrow}) p_{\downarrow}(\epsilon_{i\downarrow})$$



Two cases compared:

Spin-dependent disorder:

$$p_{\uparrow}(\epsilon) \neq p_{\downarrow}(\epsilon)$$

$$\Delta_{\uparrow} = 0; \quad \Delta_{\downarrow} \equiv \Delta$$

Spin-independent disorder:

$$p_{\uparrow}(\epsilon) = p_{\downarrow}(\epsilon)$$

$$\Delta_{\uparrow} = \Delta_{\downarrow} \equiv \Delta$$

Thermodynamic properties

Magnetization:

$$m \equiv \lim_{N_L \rightarrow \infty} \langle\langle \sum_i \hat{m}_i \rangle\rangle_{\text{dis}} / N_L$$

Double occupation:

$$d \equiv \lim_{N_L \rightarrow \infty} \langle\langle \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} \rangle\rangle_{\text{dis}} / N_L$$

Charge susceptibility:

$$\chi_c = \left(\frac{\partial n}{\partial \mu} \right)_T$$

Magnetic susceptibility:

$$\chi = \left(\frac{\partial m}{\partial h} \right)_T$$

Also others: $N_\sigma(\mu), \dots$

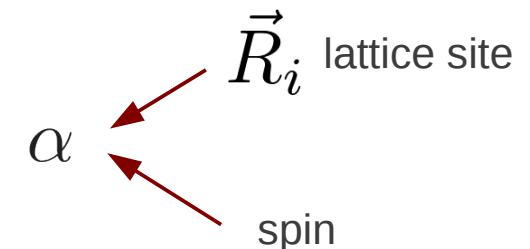
Green function

Grand canonical potential (thermodynamics):

$$\Xi = \sum_{N, \gamma_N} \langle \gamma_N | e^{-\beta(H - \mu \hat{N})} | \gamma_N \rangle$$

One-particle Green function:

$$G_{\alpha\beta}(\tau, \tau') = -\langle \langle T_\tau \hat{a}_\alpha(\tau) \hat{a}_\beta^\dagger(\tau') \rangle \rangle_{\text{dis}}$$



$$\langle \dots \rangle = \frac{1}{\Xi} \sum_{N, \gamma_N} \langle \gamma_N | e^{-\beta(H - \mu \hat{N})} \dots | \gamma_N \rangle$$

“time” ordering

$$T_\tau \hat{a}_\alpha(\tau) \hat{a}_\alpha^\dagger(\tau') = \begin{cases} \hat{a}_\alpha(\tau) \hat{a}_\alpha^\dagger(\tau') & \text{dla } \tau > \tau' \\ -\hat{a}_\alpha^\dagger(\tau') \hat{a}_\alpha(\tau) & \text{dla } \tau < \tau' \end{cases}$$

Operator in modified Heisenberg picture

$$\hat{a}_\alpha(\tau) = e^{\tau(H - \mu N)} \hat{a}_\alpha e^{-\tau(H - \mu N)}$$

Method: dynamical mean field theory

- W. Metzner, D. Vollhardt, Phys. Rev. Lett. 59, 121 (1987)
- A. Georges, G. Kotliar, W. Krauth, M. J. Rozenberg, Rev. Mod. Phys. 68, 13 (1996)
- M. Ulmke, V. Janis, D. Vollhardt, Phys. Rev. B 51, 10411 (1995)

Dyson equation:

$$G_{00}(i\omega_n) = \int_{-\infty}^{\infty} d\epsilon \frac{D(\epsilon)}{i\omega_n - \epsilon - \Sigma(i\omega_n) + \mu}$$

Semi elliptic
DOS

Consequences of high dimension limit (local problem):

$$G_{00}(\tau, \tau') = -\langle\langle c_{0\sigma}(\tau)c_{0\sigma}^*(\tau') \rangle\rangle_{S_{\text{eff}}[\mathcal{G}_0]} \text{dis}$$

$$G_{00} = \mathcal{G}_0 + \mathcal{G}_0 \Sigma G_{00}$$

QMC
Hirsch-Fye

$$S_{\text{eff}} \approx \int_0^\beta d\tau \int_0^\beta d\tau' c_{0\sigma}^*(\tau) (\mathcal{G}_0^{-1}(\tau, \tau') - \epsilon) c_{0\sigma}(\tau') + \int_0^\beta d\tau U n_{0\downarrow}(\tau) n_{0\uparrow}(\tau)$$

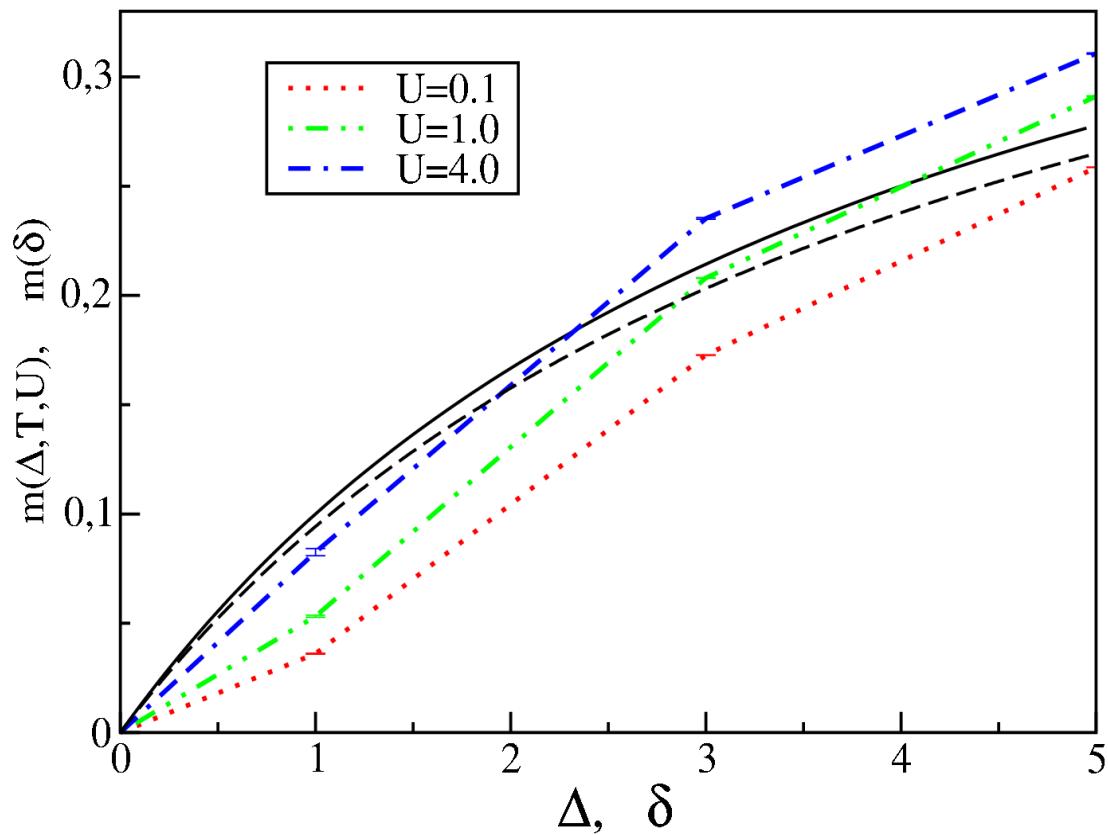
Results - magnetization

Spin-independent disorder:

$$m = 0$$

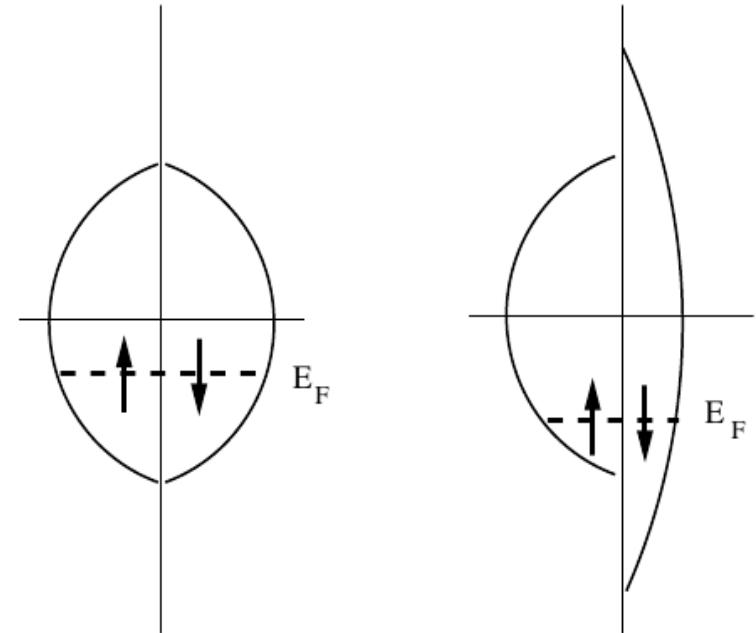
Spin-dependent disorder:

$$m \neq 0$$

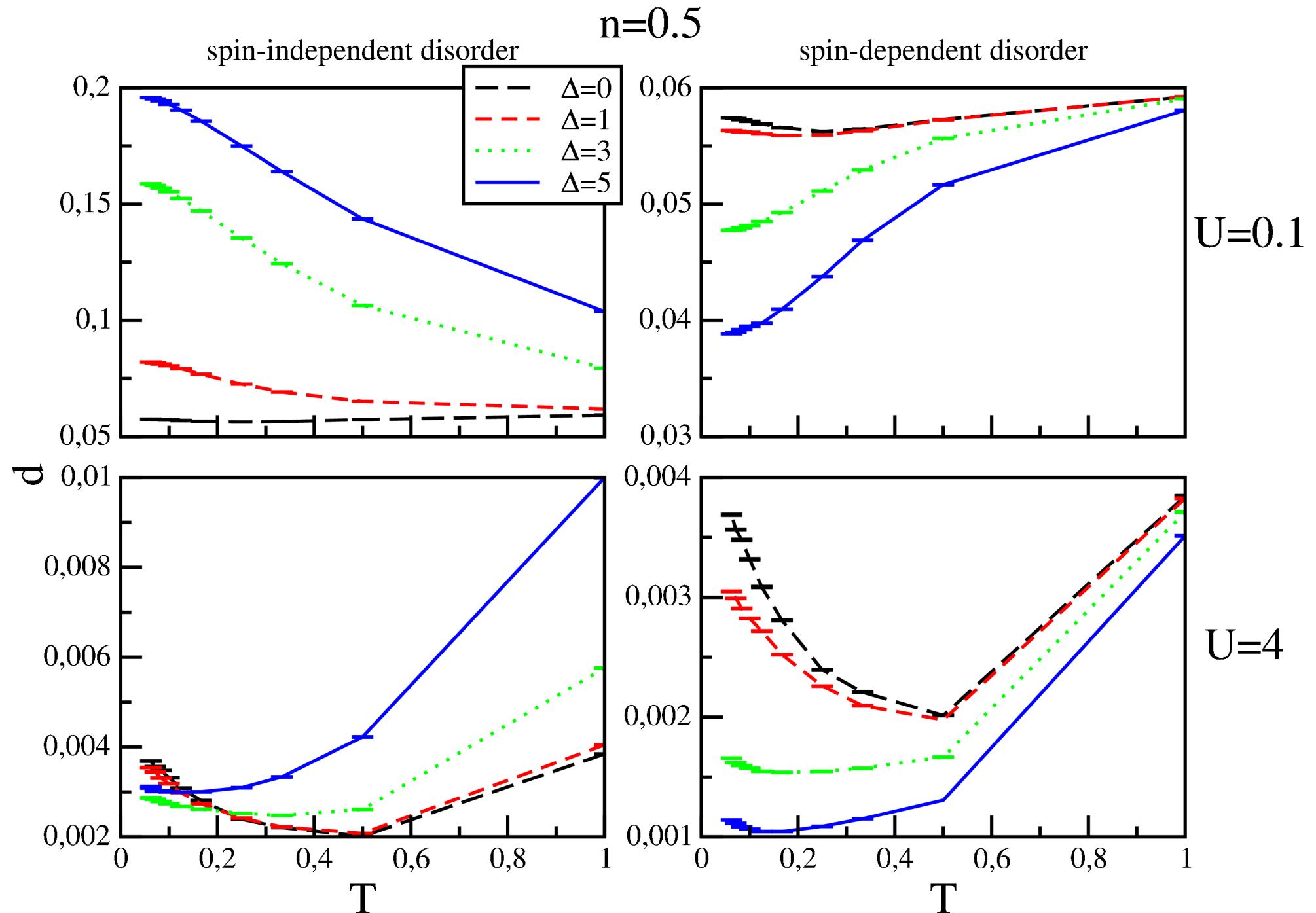


$n=0.5$; $T=1/16$

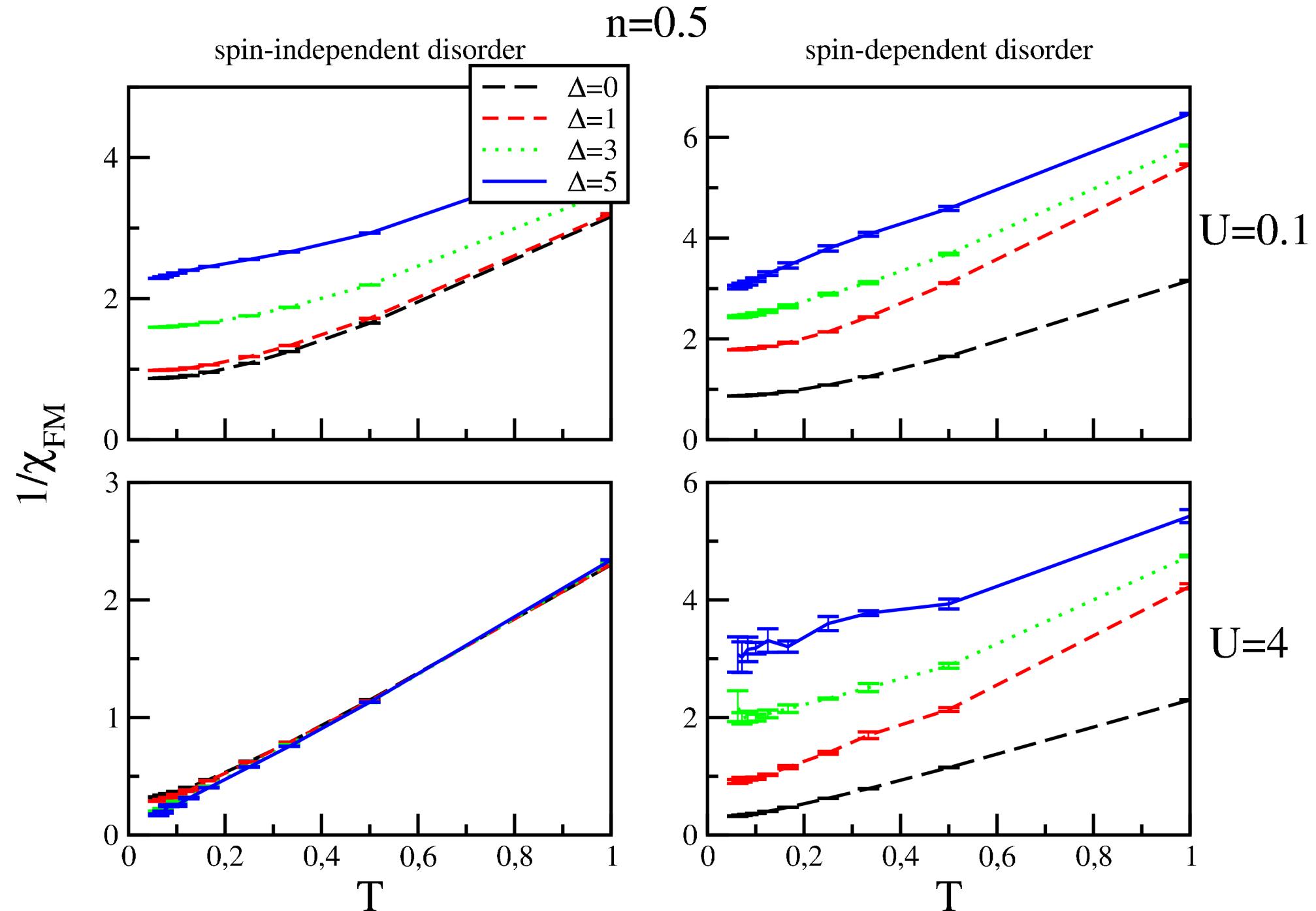
not so significant U dependence
- magnetization by noninteracting system



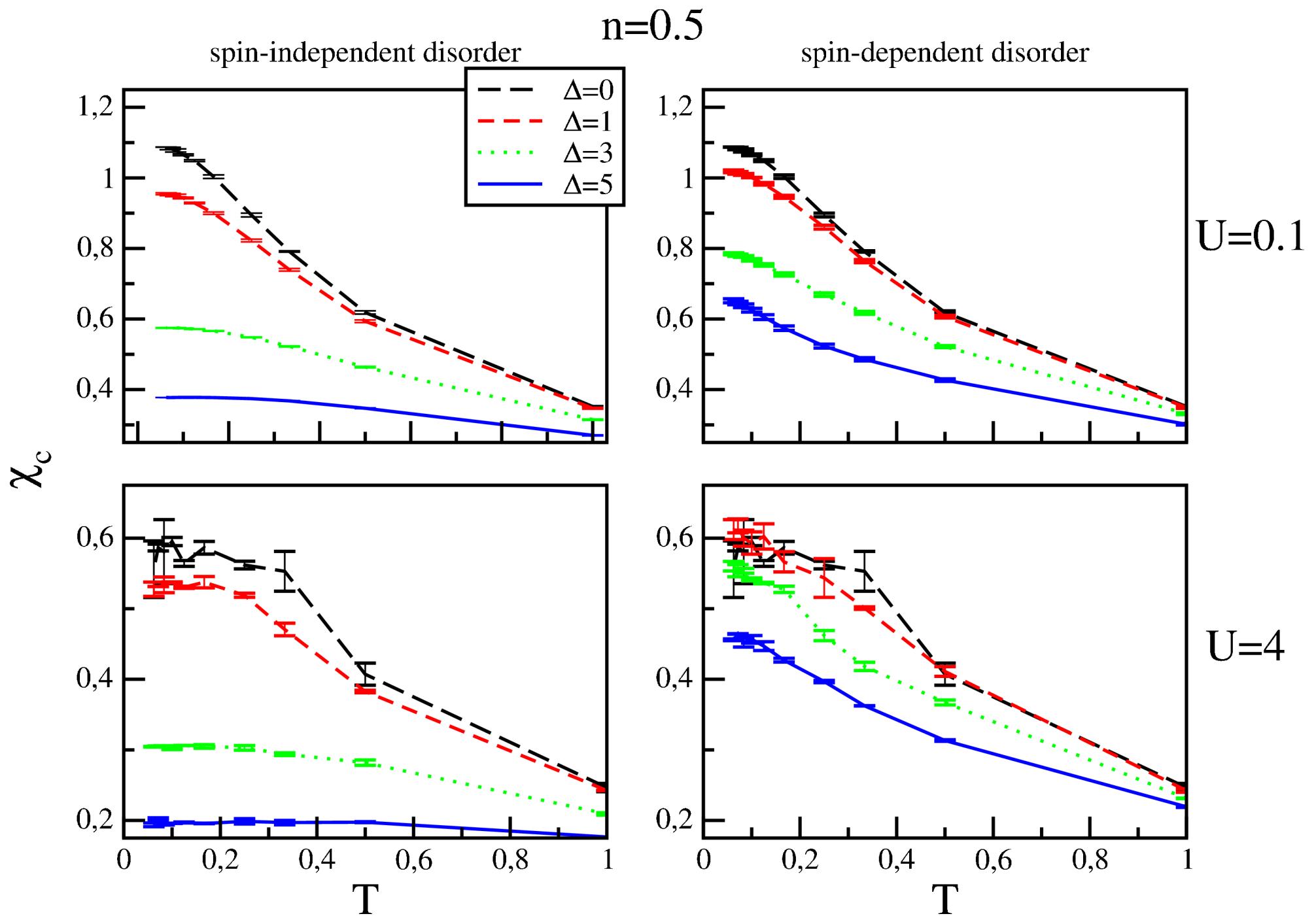
Results – double occupation



Results – ferromagnetic susceptibility



Results – charge susceptibility



Hubbard model with spin-dependent disorder

Spin imbalanced system

Spin-imbalanced system - fixed n_\uparrow, n_\downarrow

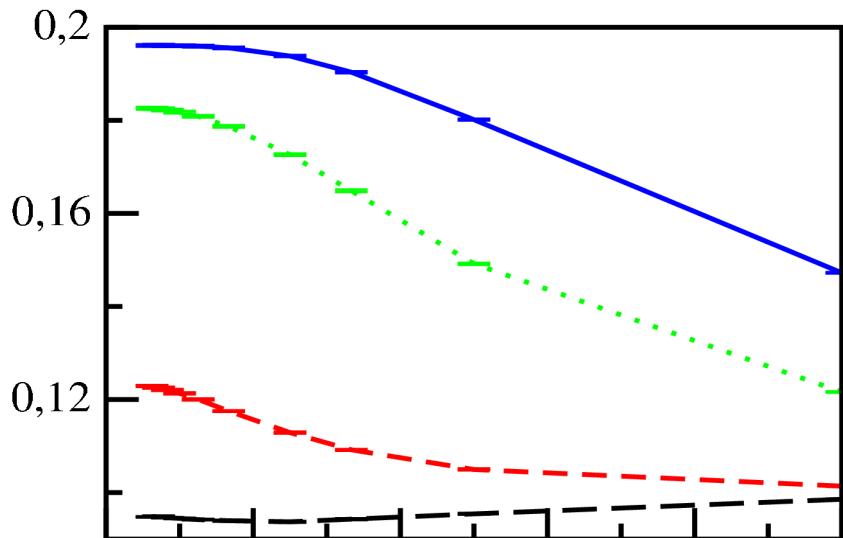
Grand canonical Hamiltonian for spin-imbalanced system:

$$H = t \sum_{(i,j)} \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} \hat{n}_{i\sigma} + \sum_i U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} - \mu_\uparrow \sum_i \hat{n}_{i\uparrow} - \mu_\downarrow \sum_i \hat{n}_{i\downarrow}$$

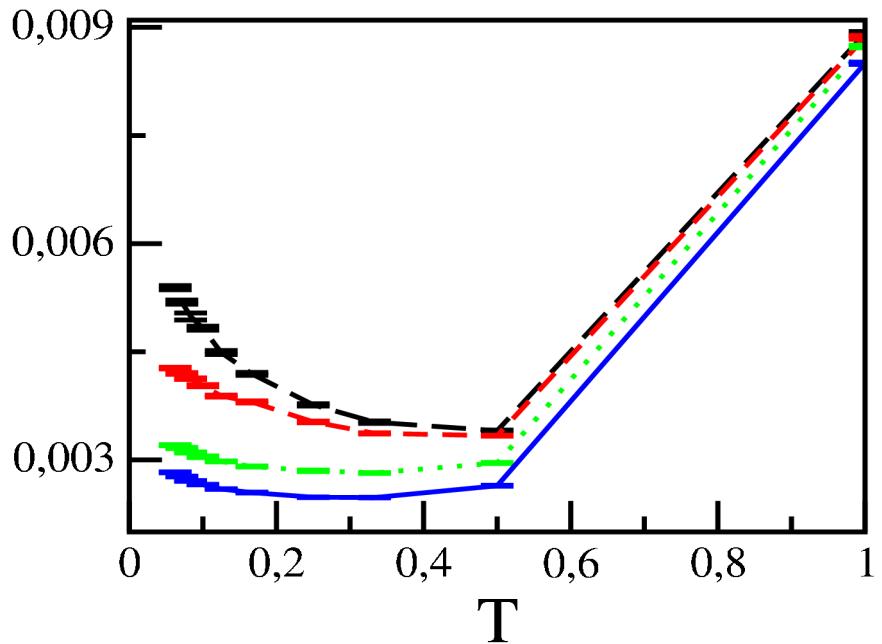
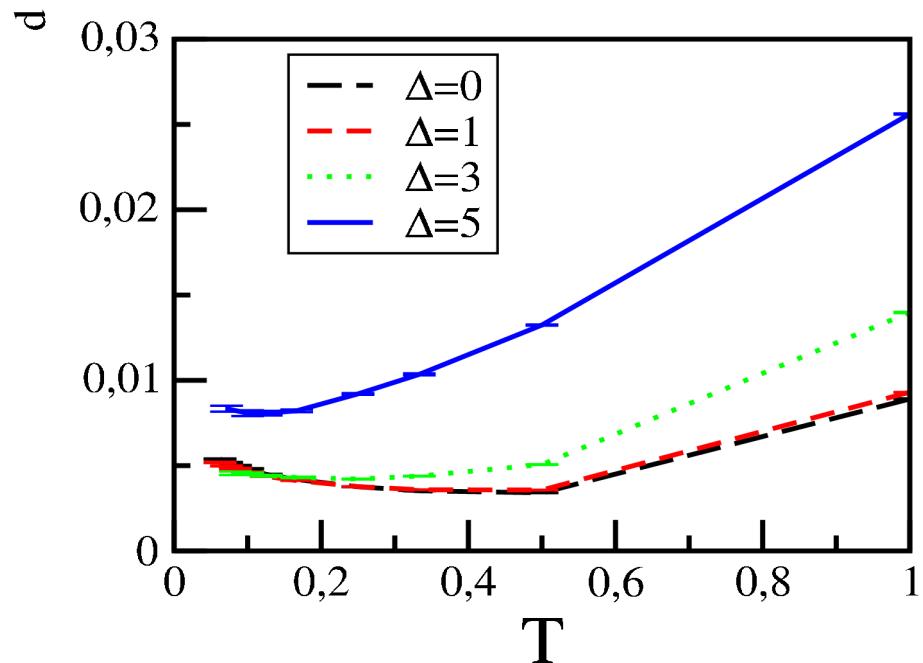
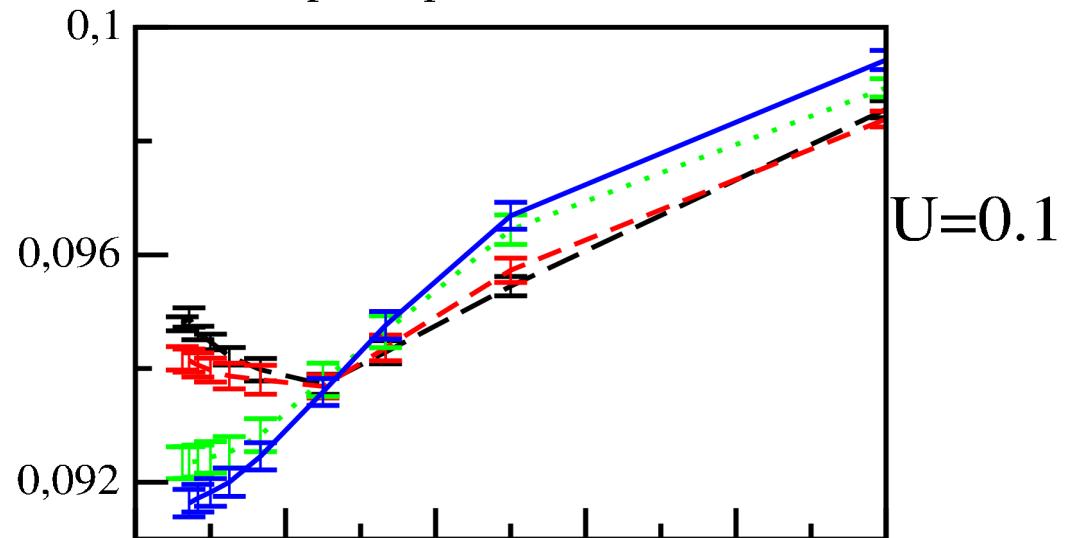
Densities fixed \rightarrow constant magnetization

Results – double occupation

$n_{\downarrow}=0.2, n_{\uparrow}=0.5$
spin-independent disorder



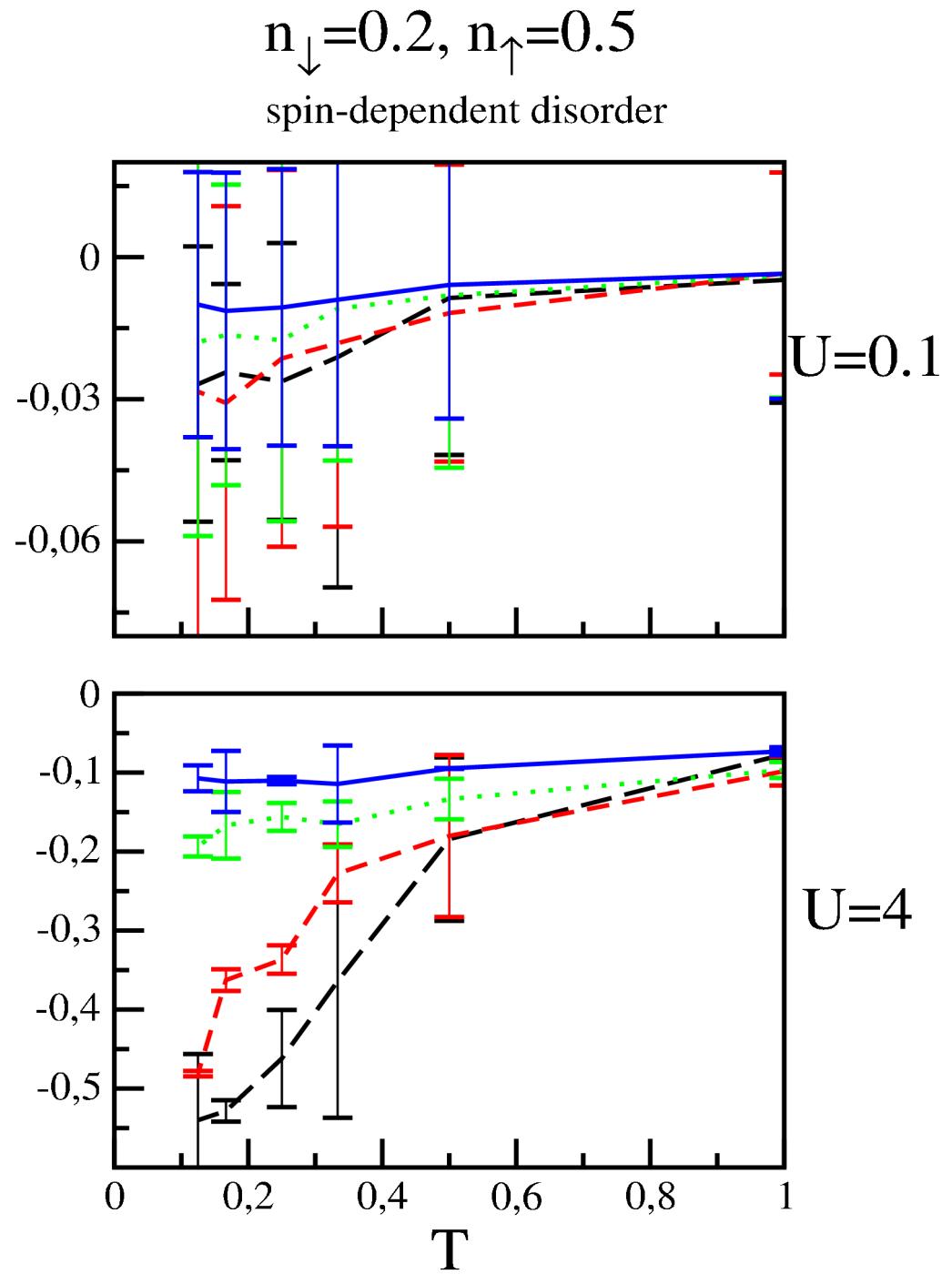
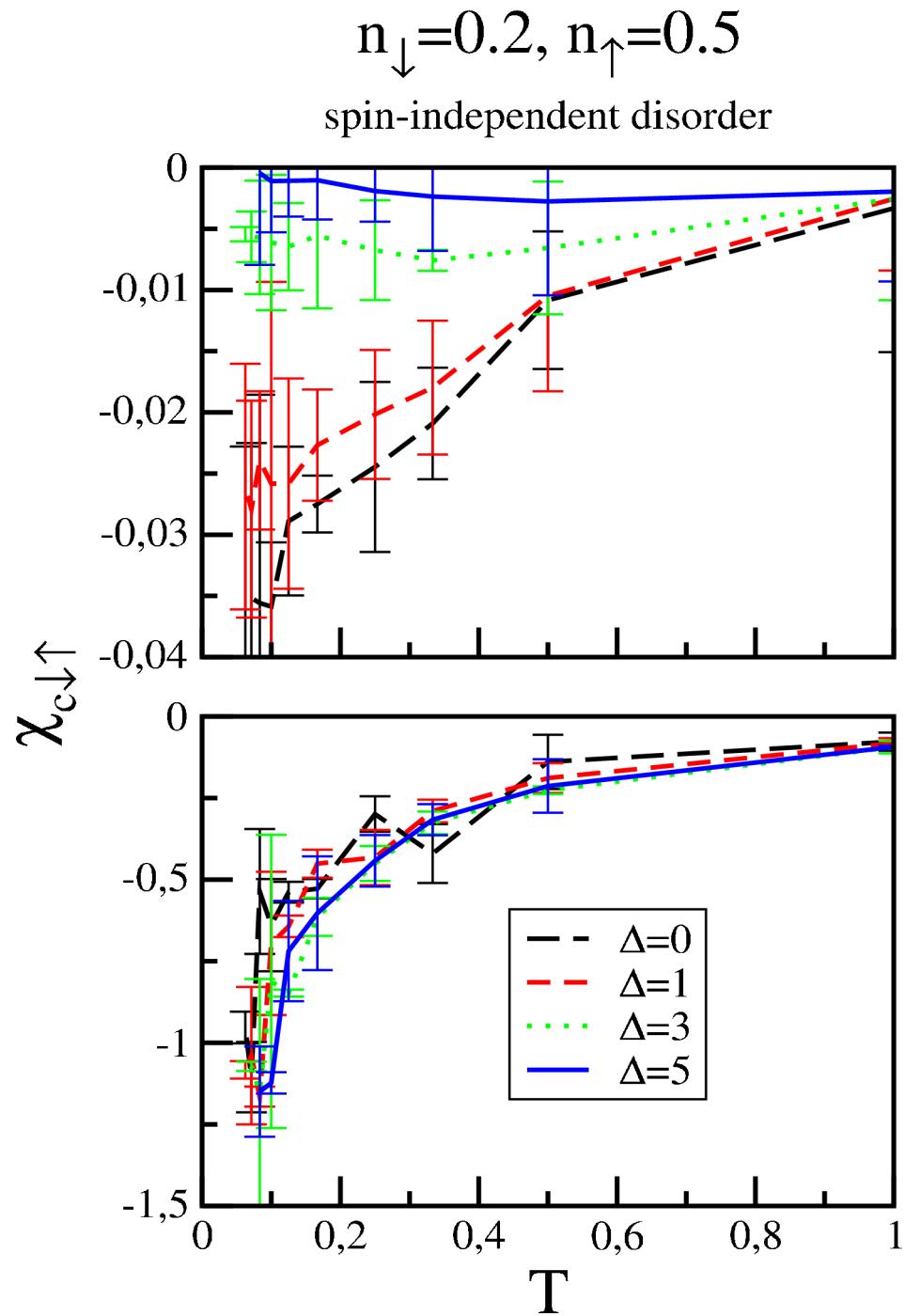
$n_{\downarrow}=0.2, n_{\uparrow}=0.5$
spin-dependent disorder



Off-diagonal compressibility

$$\chi_{c\sigma\sigma'} = \left(\frac{\partial n_\sigma(\mu_\uparrow, \mu_\downarrow, T)}{\partial \mu_{\sigma'}} \right)_T = \beta \langle \langle \hat{n}_\sigma \hat{n}_{\sigma'} \rangle - \langle \hat{n}_\sigma \rangle \langle \hat{n}_{\sigma'} \rangle \rangle_{\text{dis}} / N_L$$

Results - off-diagonal compressibility



Results – other physical properties

Other thermodynamical properties can be found in:

K. M, J. Skolimowski, P. B. Chakraborty, K. Byczuk, D. Vollhardt, 2013, arXiv:1302.3395