

Hubbard model with spin disorder

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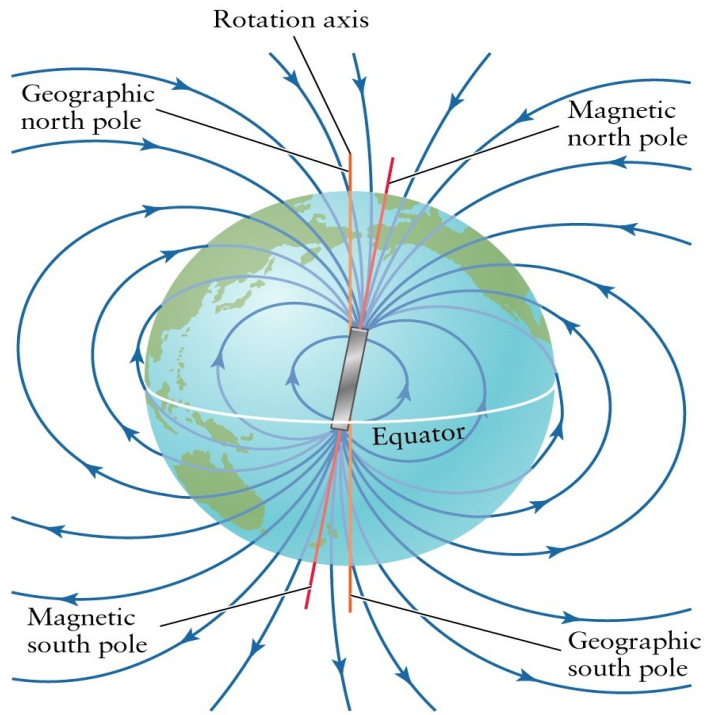
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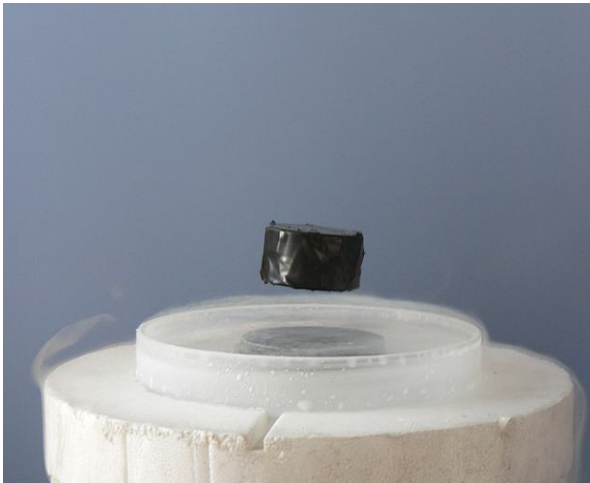
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Hubbard model



Magnetism (itinerant moment metals)



High-temperature superconductivity

Electrons in crystal

Hamiltonian

$$H = \sum_{i=1}^N \frac{(-\hbar \nabla_i)^2}{2m} + \sum_{i=1}^N V_{ions}(\vec{r}_i) + \sum_{j < i}^N \sum_{i=1}^N U_{Coulomb}(\vec{r}_i - \vec{r}_j)$$

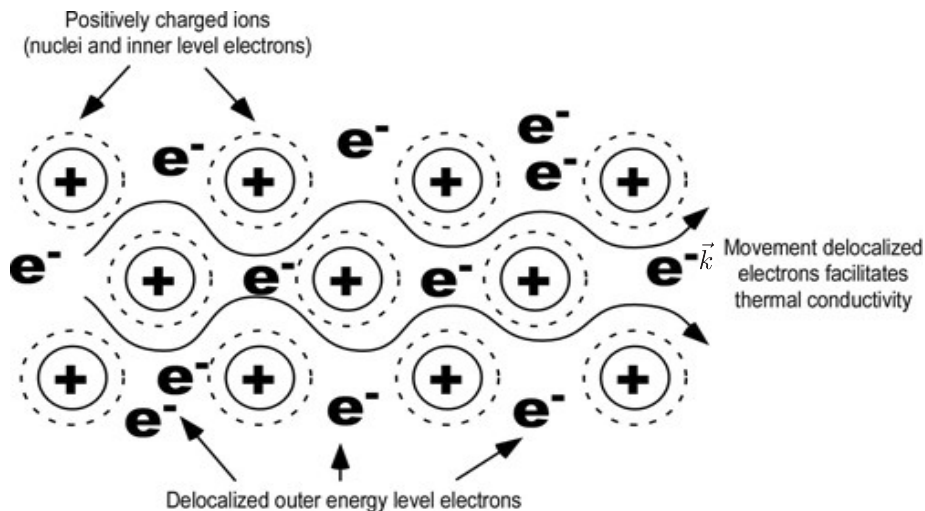
Kinetic energy

Periodic potential from ions

Coulomb interactions between electrons

Hamiltonian (second quantization)

$$H = \sum_{\alpha\beta} t_{\alpha\beta} a_{\alpha}^{\dagger} a_{\beta} + \sum_{\alpha\beta} V_{\alpha\beta} a_{\alpha}^{\dagger} a_{\beta} + \sum_{\alpha\beta\gamma\delta} a_{\alpha}^{\dagger} a_{\beta}^{\dagger} U_{\alpha\beta\gamma\delta} a_{\delta} a_{\gamma}$$



One-particle states

$n = 0, 1, \dots$

\vec{R}_i lattice site

spin

Hubbard model

From full Hamiltonian only some terms taken into consideration

-single band; local Coulomb interaction, kinetic energy term (including potential from ions):

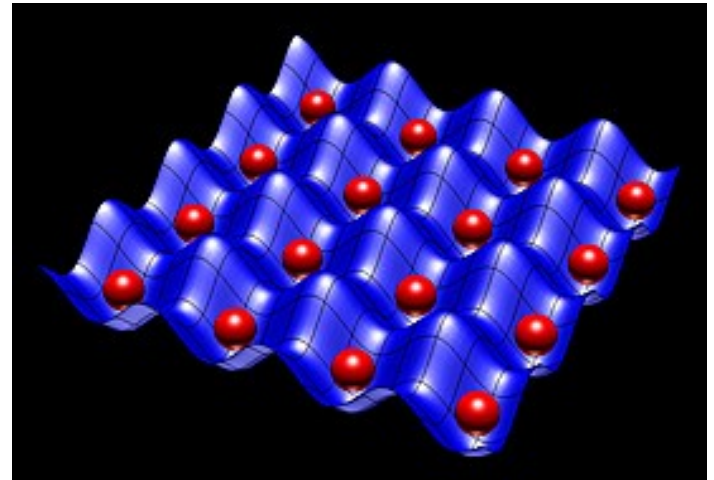
$$H = t \sum_{(i,j) \in n.n.} \sum_{\sigma=\uparrow,\downarrow} a_{i\sigma}^\dagger a_{j\sigma} + \sum_{i=1}^N U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

$$n_{i\sigma} \equiv a_{i\sigma}^\dagger a_{i\sigma}$$

Kinetic energy and
interactions with lattice

Interaction energy
between electrons

Quantum simulations of Hubbard model

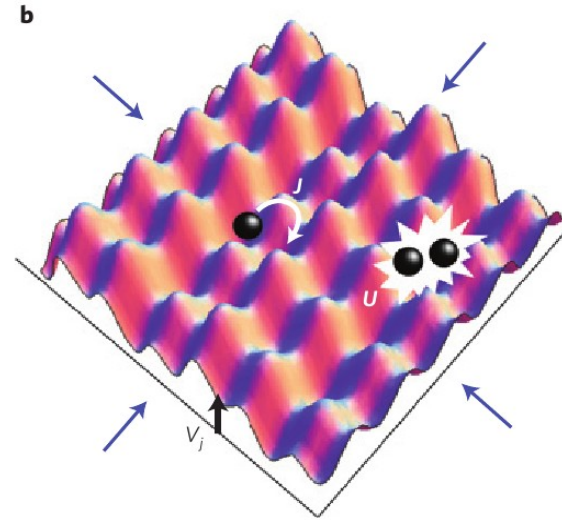


Hubbard model with disorder

$$H = t \sum_{(i,j) \in n.n.} \sum_{\sigma=\uparrow,\downarrow} a_{i\sigma}^\dagger a_{j\sigma} + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} n_{i\sigma} + \sum_{i=1}^N U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

Energy (external potential)
depends on lattice site

Quantum simulations of Hubbard model
with spin dependent disorder



How disorder influence magnetic phase?

L. Sanchez-Palencia, M. Lewenstein
Nature Physics, 2010

Type of disorder

In experiment $\epsilon_1, \dots, \epsilon_n, \dots$ fixed.

Statistical physics: averaging over probability distribution of disorder:

We assume uncorrelated rectangular disorder:

$$\mathcal{P}(\epsilon_1, \dots, \epsilon_n, \dots) = \prod_i P(\epsilon_i)$$

$$P(\epsilon) = \begin{cases} \frac{1}{\Delta} & |\epsilon| < \frac{\Delta}{2} \\ 0 & |\epsilon| > \frac{\Delta}{2}. \end{cases}$$

What we investigate?

How interplay of spin disorder and interactions influence

-phase (paramagnetic, antiferromagnetic)

-susceptibility

-conductivity

-compressibility

Grand canonical potential (thermodynamics):

$$\Xi = \sum_{N, \gamma_N} \langle \gamma_N | e^{-\beta(H - \mu \hat{N})} | \gamma_N \rangle$$

One-particle Green function:

$$G_{\alpha\beta}(\tau, \tau') = -\langle T_\tau \hat{a}_\alpha(\tau) \hat{a}_\beta^\dagger(\tau') \rangle$$

$$\langle \dots \rangle = \frac{1}{\Xi} \sum_{N, \gamma_N} \langle \gamma_N | e^{-\beta(H - \mu \hat{N})} \dots | \gamma_N \rangle$$

“time” ordering

$$T_\tau \hat{a}_\alpha(\tau) \hat{a}_\alpha^\dagger(\tau') = \begin{cases} \hat{a}_\alpha(\tau) \hat{a}_\alpha^\dagger(\tau') & \text{dla } \tau > \tau' \\ -\hat{a}_\alpha^\dagger(\tau') \hat{a}_\alpha(\tau) & \text{dla } \tau < \tau' \end{cases}$$

Operator in Heisenberg picture

$$\hat{a}_\alpha(\tau) = e^{\tau(H - \mu N)} \hat{a}_\alpha e^{-\tau(H - \mu N)}$$

Green function for system with perturbation

$$H = H_0 + H_1$$

For H_0 :
$$G_{0,\alpha\beta}(\tau, \tau') = -\langle T_\tau \hat{a}_\alpha(\tau) \hat{a}_\beta^\dagger(\tau') \rangle$$

Hamiltonian in “average” and in „time” evolution

For H :
$$G_{\alpha\beta}(\tau, \tau') = -\langle T_\tau \hat{a}_\alpha(\tau) \hat{a}_\beta^\dagger(\tau') \rangle$$

Dyson equation:

$$G_0^{-1} = G^{-1} + \Sigma$$

Self-energy (diagrams)

Noninteracting system with disorder

$$H = t \sum_{(i,j) \in n.n.} \sum_{\sigma=\uparrow,\downarrow} a_{i\sigma}^\dagger a_{j\sigma} + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} n_{i\sigma}$$

Green function without disorder: $[G_0(\omega)]_{i\sigma j\sigma}$

Green function for system with disorder (Dyson equation):

$$[G_0^{-1}(\omega)]_{i\sigma j\sigma'} = [G^{-1}(\omega)]_{i\sigma j\sigma'} + \delta_{ij} \delta_{\sigma\sigma'} \epsilon_i$$

$$G_0^{-1} = G^{-1} + \epsilon$$

Green function averaged over disorder:

$$\langle G \rangle_{dis}$$

Coherent potential approximation

Dyson equation:

$$G_0^{-1} = G^{-1} + \epsilon$$

Definition of coherent potential Σ , which the same as ϵ is diagonal: $[\Sigma]_{ij} = \delta_{ij}\Sigma$

$$G_0^{-1} = \langle G \rangle^{-1} + \Sigma$$

yields:

$$G = \langle G \rangle + \langle G \rangle T \langle G \rangle$$

$$\text{where } T = (\epsilon - \Sigma)[1 - \langle G \rangle(\epsilon - \Sigma)]^{-1}$$

$$\text{Average leads to: } \langle T \rangle = 0$$

Coherent potential approximation:
leads to closed system of equations.

$$[\langle G \rangle]_{ij} \approx \delta_{ij}[\langle G \rangle]_{ii}$$

Hubbard model with disorder

$$H = t \sum_{(i,j) \in n.n.} \sum_{\sigma=\uparrow,\downarrow} a_{i\sigma}^\dagger a_{j\sigma} + \sum_i \sum_{\sigma=\uparrow,\downarrow} \epsilon_{i\sigma} n_{i\sigma} + \sum_{i=1}^N U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

Dynamical mean field approach

Metzner i Vollhardt, 1987 – limit of high dimensional lattice (Georges et al. (1996))

Dynamical mean field equations for disorder

Dyson equation:

$$G = G_0 + G_0 \Sigma G$$

Consequences of high dimension limit (local problem):

$$G_{00}(\tau, \tau') = -\langle \langle T_\tau c_{0\sigma}(\tau) c_{0\sigma}^*(\tau') \rangle \rangle_{S_{\text{eff}}[\mathcal{G}_0]} \text{dis}$$

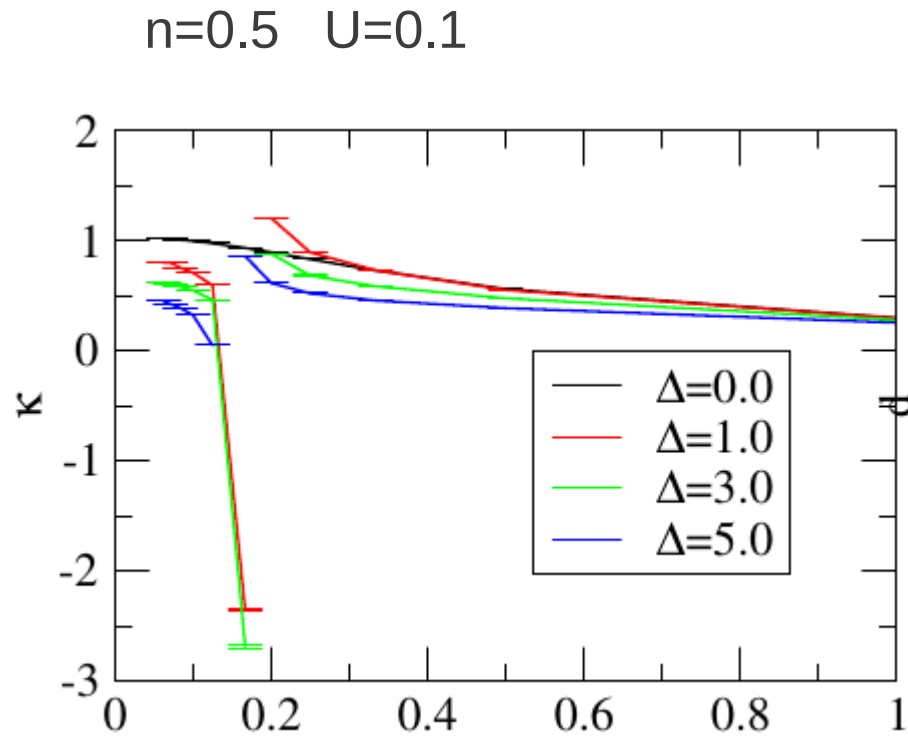
← QMC

$$G_{00} = \mathcal{G}_0 + \mathcal{G}_0 \Sigma G_{00}$$

$$S_{\text{eff}} \approx \int_0^\beta d\tau \int_0^\beta d\tau' c_{0\sigma}^*(\tau) (\mathcal{G}_0^{-1}(\tau, \tau') - \epsilon_0) c_{0\sigma}(\tau') + \int_0^\beta d\tau U n_{0\downarrow} n_{0\uparrow}$$

Actual results

Without disorder ($\Delta=0$) – proper results.
Problem for disorder ($\Delta>0$)



Thermodynamics: compressibility > 0

Where is the problem?

Fourier transform?

Asymptotic of Green function:

$$G(\omega) \sim 1/\omega \quad \text{for } \omega \rightarrow \infty$$

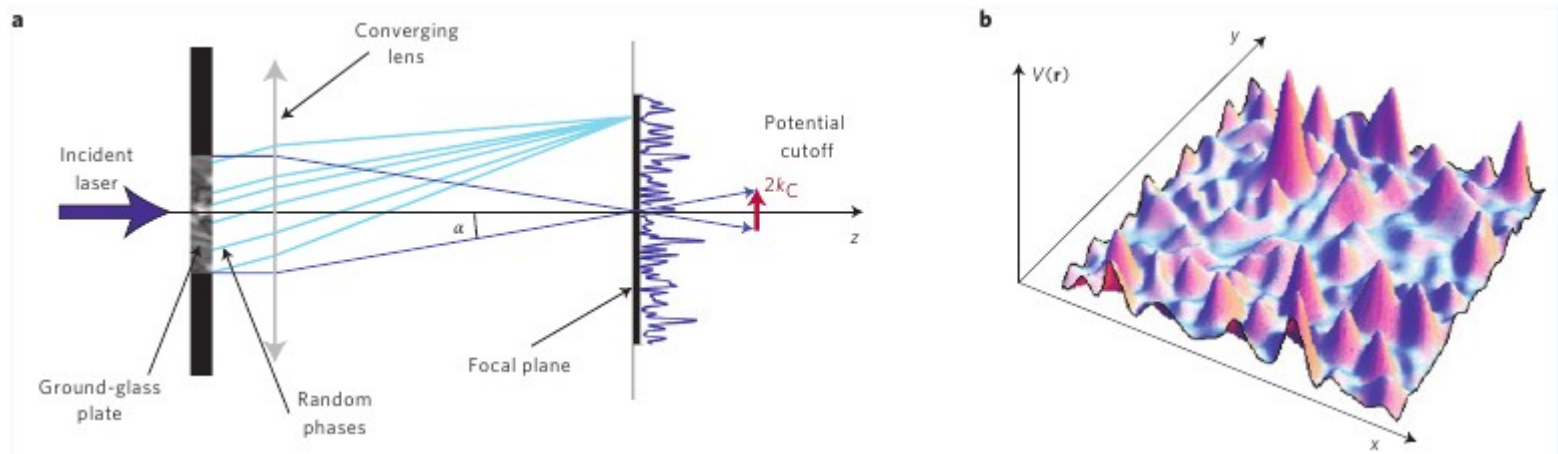
(Fourier transforms needs to be calculated properly)

No (program improved –
Ulmke trick avoided)

Transport coefficients?

To be checked...

Disorder in cold atoms



L. Sanchez-Palencia, M. Lewenstein
Nature Physics, 2010

Compressibility:

$$\kappa = \frac{1}{n^2} \left(\frac{\partial n}{\partial \mu} \right)_T \quad \kappa^{-1} \equiv -V \left(\frac{\partial p}{\partial V} \right)_T$$

Charge susceptibility:

$$\chi_c = \left(\frac{\partial n}{\partial \mu} \right)_T$$

Magnetic susceptibility:

$$\chi = \left(\frac{\partial m}{\partial h} \right)_T$$

Disorder and interaction competition

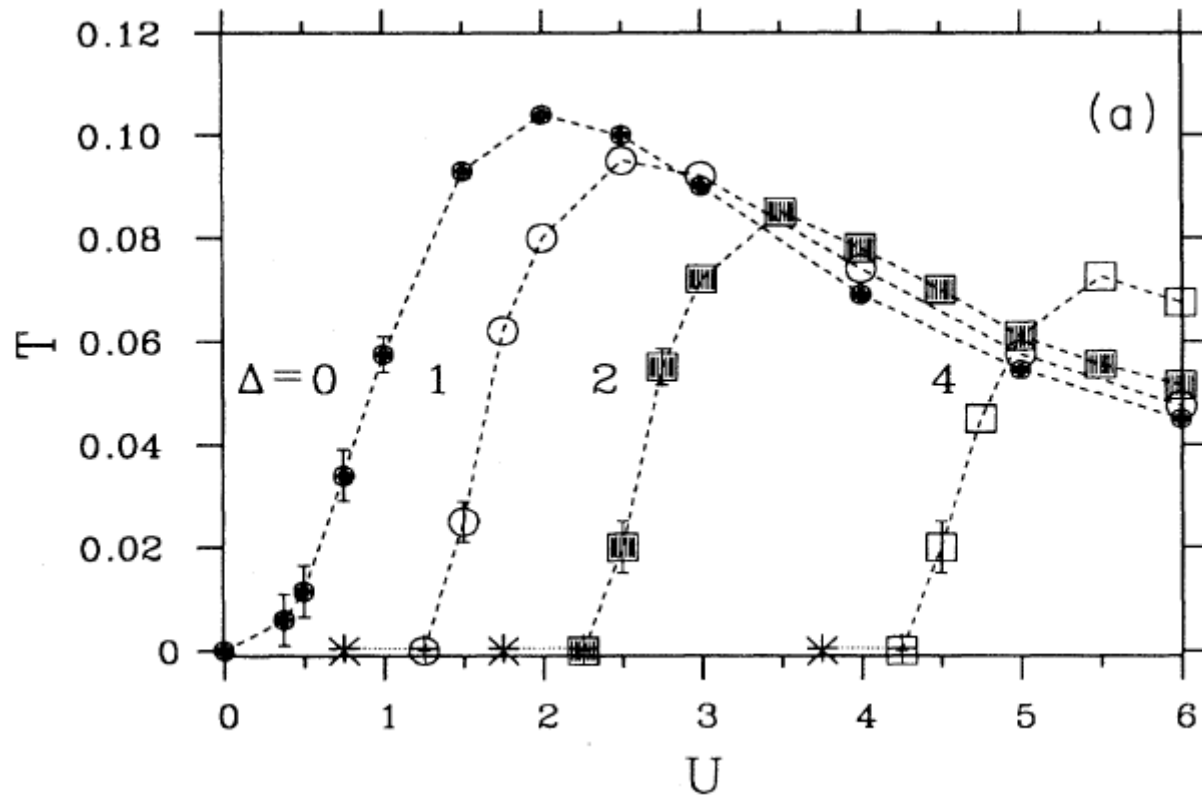


FIG. 4. (a) $T - U$ -phase diagram for the binary alloy with $\Delta = 0, 1, 2, 4$ obtained from the zeroes of χ_{AF}^{-1} (see Fig. 3). The AF phase is stable below the curves. The dotted lines at $T = 0$ depict the regimes where the Curie law would give negative transition temperatures. Below the crosses χ_{AF}^{-1} has no zeroes but a minimum and an AF phase can no longer be expected. (b) $T - U$ -phase diagram for the semielliptic distribution of the random energies with width $\Delta = 0, 2, 4, 6$.

Hubbard model

Georges et al. (1996): cuprate (tlenki miedzi) superconductors (transition metals – niecałkowicie zapełniona powłoka d)

Ulmke trick

It smoothed Green function(?)

$$\bar{G}_n = \delta\tau / [1 - \exp(\delta\tau/G_n)],$$

$$\bar{G}(\tau) = \beta^{-1} \sum_n e^{i\omega_n\tau} \bar{G}_n,$$

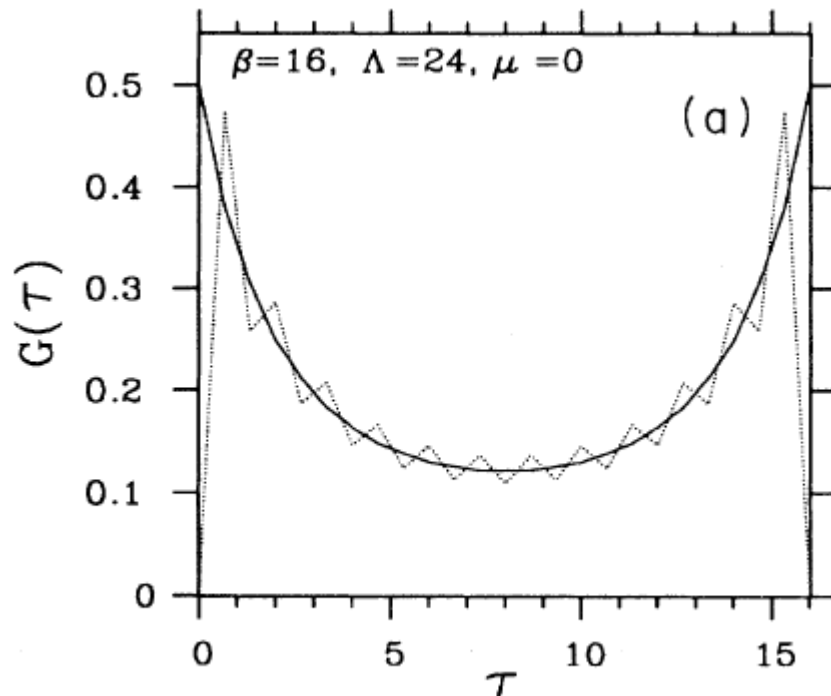


FIG. 1. Local Green function $G(\tau)$ for $U = 0$ and $\mu = 0$ (a), $\mu = 0.5$ (b) obtained by usual Fourier transformation (dotted line) and by redefinition according to Eq. (13) (solid line).

Picture from: M. Ulmke, V. Janis, D. Vollhardt (1995)